

Rule for Mode Coupling Efficiency in Optical Waveguide Crossing

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Abstract – Crossing an optical waveguide requires a beam coupling from free space to waveguide at the entrance plane and another beam coupling from waveguide to free space at the exit plane of the waveguide. The aim of this paper is to provide a simple rule expressing the relationship between the involved numbers of free and guided modes that efficiently rebuild the field at each end of the waveguide. Using a numerical program built on Maple software, the rule was determined to be effective independently of the ratio between the beam spot size and the waveguide radius.

Index Terms – Electromagnetic field, mode coupling, optical waveguide.

I. INTRODUCTION

Hollow dielectric optical waveguides are widely used in optical systems such as resonators [1, 2], optical transmission and communication systems [3], circular waveguide filters [4], and, nowadays, in integrated optics [5].

Many authors have described light propagation in cylindrical optical waveguides [6–9]. Transverse modes inside a waveguide can be divided into three families: transverse electric (TE), transverse magnetic (TM), and hybrid (EH) modes. Each family constitutes a complete and orthogonal set for expressing the radially symmetric electromagnetic field in the waveguide.

Apart from the propagation inside the waveguide, two check points need to be considered when crossing a waveguide which are its entrance and exit ports. At the entrance port, mode coupling occurs from free space to confined one and vice versa at the exit port. For both ends of the waveguide, mode coupling has been studied by considering the ratio between the beam spot size w and the radius of the waveguide a [10–13].

At the waveguide entrance plane, Smith [10] described earliest experiments performed by matching

the fundamental mode from a conventional He–Ne laser into a hollow dielectric waveguide. The incident beam was focused in such a way that the beam waist w_0 occurred at the entrance plane of the dielectric waveguide. It is worth noting that Smith performed transmission measurement by matching the fundamental free mode into the fundamental guided mode. For a good matching, he experimentally found a ratio $w_0/a = 0.49$ which was however considerably different from the theoretical ratio $w_0/a = 0.728$ he considered.

Always for the entrance plane and as a function of w/a , Roullard and Bas [11] gave the fraction of the coupled energy from the free fundamental Gaussian mode TEM_{00} into the two first waveguide modes. They mentioned that an efficient coupling happens at $w/a = 0.502$. This value is different from the value $w/a = 0.6435$ that maximizes coupling with only the first-order waveguide mode as given by Abrams and Chester [12] and Tack [13].

At the waveguide exit plane, the useful approach to release mode coupling is founded on the use of a small number of free space modes but with a choice of a specific ratio w/a . Indeed, as mentioned by Guerlach [14], to minimize the truncation error in expanding the field emerging from the waveguide, a small value of the ratio is favored. But a small w results in a large divergence of the beam field and consequently a large truncation error. Avoiding this constraint of the w/a ratio choice requires a considerable number of modes with numerical problem consequences.

To conclude, either at the entrance or at the exit plane of the waveguide, mode coupling using a reduced number of modes is conditioned by the choice of the appropriate ratio w/a . Working with an arbitrary ratio, requires a large number of modes that unfortunately leads to numerical problems. Therefore, the question that we ask here is about the optimum number of modes that we should consider for realizing an efficient mode

coupling regardless of a specific ratio w/a . To reach this goal, a simple rule is then given in this paper.

II. THEORY AND NUMERICAL RESULTS

Let us consider a hollow dielectric waveguide with a circular cross section of radius a , a length L , and a complex refractive index v . Assume that an incident beam is coming in from the left side as shown in Figure 1.

Crossing an optical waveguide is carried out in three successive steps which are: mode coupling at the entrance plane, propagating the derived guided electromagnetic field along the length L inside the waveguide, and finally performing mode coupling again at the exit plane to go back to free space. Figure 2 shows the sequence of operations undertaken in the waveguide crossing. This process requires normalized functions for free space and waveguide modes to be given.

In free space for the case of cylindrical symmetry, the normalized Laguerre Gauss functions are given by [15]

$$\begin{aligned} \text{TEM}_{pl} &= \left(\frac{2}{w\sqrt{1+\delta_{0l}}} \sqrt{\frac{m!}{\pi(l+m)!}} \right) \left(\frac{r}{w\sqrt{2}} \right)^l L_m^l \left(2\frac{r^2}{w^2} \right) \\ &\times \exp \left\{ -\frac{r^2}{w^2} - j\frac{kr^2}{2R} \right\} \begin{cases} \cos(l\varphi) \\ \sin(l\varphi) \end{cases}, \end{aligned} \quad (1)$$

where r , φ are the cylindrical coordinates, L_m^l is the generalized Laguerre polynomials, δ is the Kronecker symbol, j is the complex number such as $j^2 = -1$, and k is the wave number; R and w are the phase front radius and the spot size, respectively.

For the waveguide, the normalized hybrid functions with circular symmetry field are [14, 16]

$$\text{EH}_{1n}(r) = \begin{cases} \frac{1}{\sqrt{\pi|J_1(u_{1n})|a}} J_0\left(\frac{u_{1n}r}{a}\right) & r \leq a \\ 0 & r > a \end{cases}, \quad (2)$$

where J_0 and J_1 are the Bessel functions and u_{1n} is the n th root of the former one.

Mode coupling at each end of the waveguide is expressed using coupling coefficients C_{mn} . These coefficients are reached by equalizing the electromagnetic fields at both sides of the considered waveguide port and

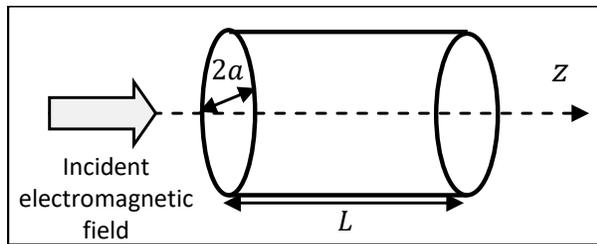


Fig. 1. Incident beam to a cylindrical optical waveguide.

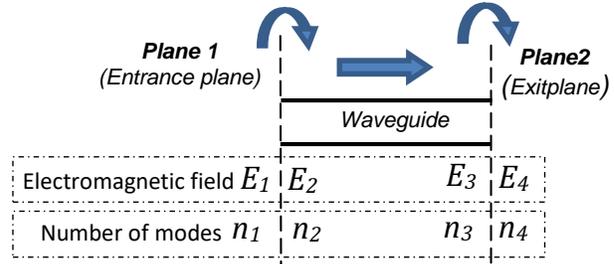


Fig. 2. Sequence of operations undertaken in waveguide crossing.

they are given by [14]

$$C_{mn} = 2\pi \int_0^a r \text{TEM}_m(r) \text{EH}_n(r) dr. \quad (3)$$

As shown in Figure 2, mode coupling occurs at plane 1 using a number of excited guided modes n_2 and at plane 2 using a number of excited free space modes n_4 . Simple rules giving the optimum n_2 and n_4 numbers for an efficient mode coupling are then the subject of Sections II-A and II-C.

A. Rule for mode coupling at the waveguide entrance plane

The field equality $E_1 = E_2$ inside the waveguide opening, at the entrance plane, is required for an efficient mode coupling. At both sides of this plane, the electromagnetic field is

$$E_1(r) = \sum_{m=1}^{n_1} a_m \text{TEM}_m(r), \quad (4)$$

for the free space and

$$E_2(r) = \sum_{n=1}^{n_2} C_n \text{EH}_n(r), \quad (5)$$

for the waveguide.

The integers n_1 and n_2 are the number of free space and waveguide modes, respectively; a_m represents the expansion coefficients of the incoming free field, whereas C_n are the derived expansion coefficients of the calculated guided field using eqn (3):

$$C_n = 2\pi \sum_{m=1}^{n_1} a_m C_{mn}. \quad (6)$$

Table 1 gives, for each given number n_1 of the used free space modes, the optimum number n_2 of waveguide modes which allows an efficient mode coupling at the waveguide entrance plane. This result is achieved for different w/a ratios.

The plot n_2 versus n_1 is illustrated by the point curve in Figure 3.

Applying the mathematical fit instructions of Maple software using the data from Table 1, the appropriate function $n_2 = f(n_1)$, shown in solid line in Figure 3, is deduced and the following rule for mode coupling at the

Table 1: Optimized number n_2 of waveguide modes as a function of the number n_1 of the used free space modes

n_1	n_2
3	7
8	8
11	9
15	9
19	10
25	12
38	13
45	15
60	17
80	19
100	21

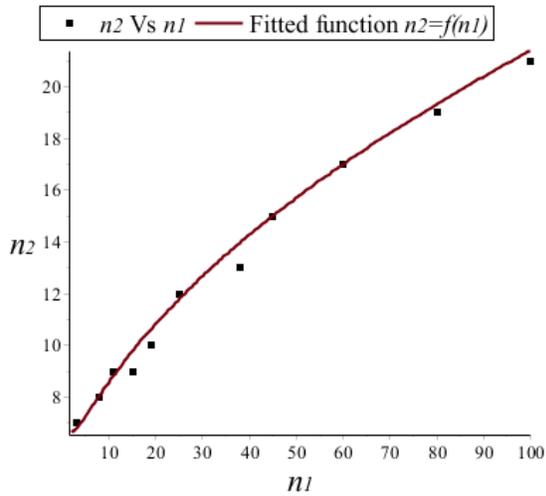


Fig. 3. Optimized number n_2 of waveguide modes as a function of the number n_1 of the used free space modes.

waveguide entrance plane is obtained:

$$n_2 = \left\lfloor 2 \left(\sqrt{n_1} + \frac{2}{\sqrt{n_1}} \right) + 1 \right\rfloor, \quad (7)$$

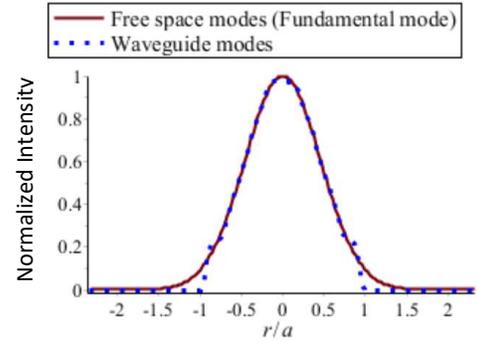
where the notation $\lfloor \square \rfloor$ is the floor function.

For different n_1 and according to this rule, Figure 4 shows a very good agreement in field superposition for mode coupling at the entrance plane of the waveguide.

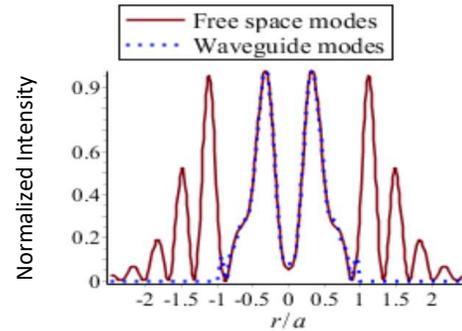
The results shown in Figure 4 (a)–(c) have been computed with $w/a = 0.91$, $w/a = 0.8$, and $w/a = 0.66$, respectively. But we emphasize that for each case, and thanks to the rule, an efficient mode coupling has been confirmed for other arbitrary ratios w/a .

B. Propagation inside the waveguide

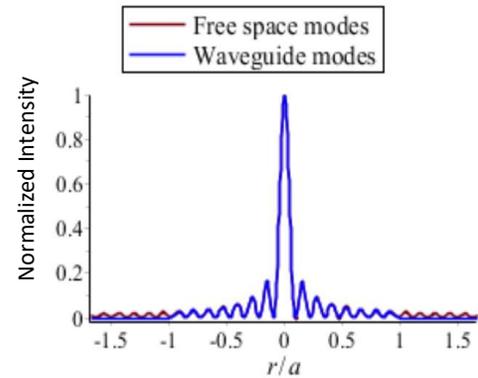
Attenuation and relative phase shift of the propagating field, over a distance l through the waveguide, are calculated by multiplying each guided mode by



(a)



(b)



(c)

Fig. 4. Superposition of transverse intensity patterns at the waveguide entrance plane with $a = 0.6$ mm for (a) $n_1 = 1$, (b) $n_1 = 13$, and (c) $n_1 = 35$.

the corresponding element of the following diagonal matrix [14]:

$$CG_{mn} = \delta_{mn} \exp \left[u_{1m}^2 \left(\frac{l}{ka^2} \right) \operatorname{Re} \left(\frac{v_n}{ka} \right) \right] \times \exp \left[-\frac{i}{2} (u_{1m}^2 - u_{11}^2) \left(\frac{l}{ka^2} \right) \left(1 + 2 \operatorname{Im} \left(\frac{v_n}{ka} \right) \right) \right], \quad (8)$$

where δ_{mn} is the Kronecker symbol, $k = 2\pi/\lambda$, u_{1m} is the m th zero of J_0 and $v_n = [v^2+1]/2[v^2-1]^{\frac{1}{2}}$ where v is the complex refractive index of the waveguide material.

Figure 5 shows the electrical field intensity in different positions inside a cylindrical optical waveguide.

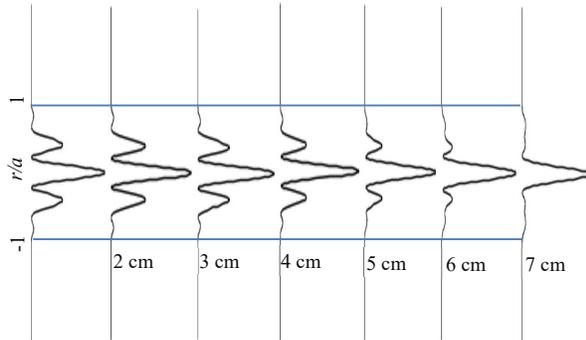


Fig. 5. Transverse intensity patterns at different positions inside an optical waveguide of length $L = 7$ cm and $a = 0.6$ mm.

C. Rule for mode coupling at the waveguide exit plane

Here, the set of waveguide modes must be coupled to a set of free space modes. The field at the left side of the waveguide exit plane is

$$E_3(r) = \sum_{n=1}^{n_3} CG_n EH_n(r), \quad (9)$$

where the integer n_3 is the number of the waveguide modes, which is equal to n_2 obtained from eqn (7), and the CG_n coefficients are calculated using eqn (8).

The field at the right side of the waveguide exit plane is

$$E_4(r) = \sum_{m=1}^{n_4} C_m TEM_m(r), \quad (10)$$

where C_m are defined in eqn (6) and the integer n_4 is the considered number of the TEM modes.

Table 2 shows, for each given number n_3 of the used waveguide modes, the optimum number n_4 of free space modes which allows an efficient mode coupling at the waveguide exit plane. This result is achieved for different w/a ratios.

The plot n_4 versus n_3 is illustrated by the point curve in Figure 6.

Applying the mathematical fit instructions of Maple software, using the data from Table 2, the appropriate function $n_4=f(n_3)$ is deduced and the following rule for mode coupling at the waveguide exit plane is obtained:

$$n_4 = \lfloor 0.35n_3^2 + 1 \rfloor. \quad (11)$$

Table 2: Optimized number n_4 of free space modes as a function of the number n_3 of the used waveguide modes

n_3	n_4
7	16
8	21
9	29
10	36
11	40
12	55
14	70
17	100
19	130
22	175
25	219
30	316

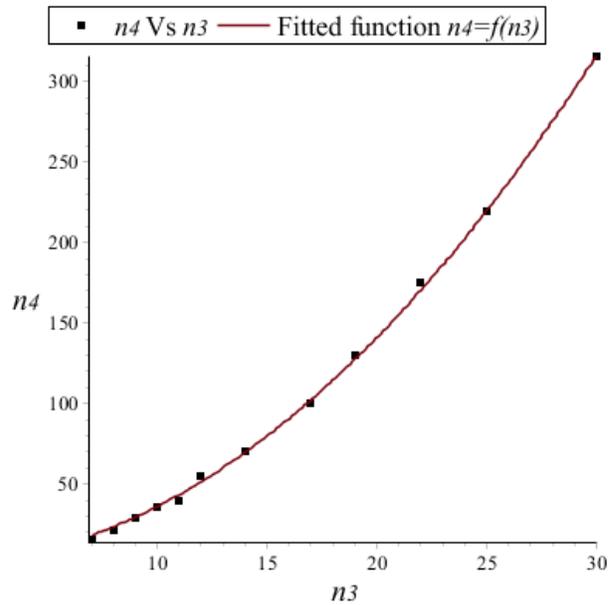


Fig. 6. Optimized number n_4 of free space modes as a function of the number n_3 of the used waveguide modes.

Figure 7 shows mode coupling according to this rule for different values of n_3 .

The results shown in Figure 7 (a)–(c) have been computed with $w/a = 0.58$, $w/a = 0.5$, and $w/a = 0.8$, respectively. As mentioned in Section II-A, we confirm that for each case, and thanks to the rule, an efficient mode coupling has been verified for other arbitrary ratios w/a .

It should be noted that at both ports of the waveguide and according to the rule, mode coupling goes successfully independent of the ratio w/a for values outside the range $0.45 < w/a < 0.73$ applied by different authors [10–14]. Indeed, the rule works very well

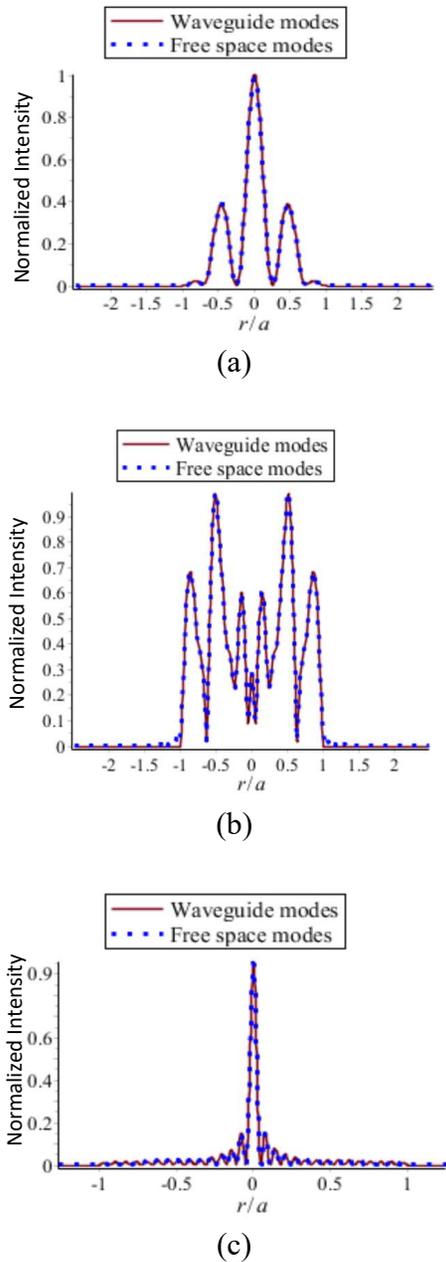


Fig. 7. Superposition of transverse field distribution at the waveguide exit plane with $a = 0.6$ mm for (a) $n_3 = 7$, (b) $n_3 = 10$, and (c) $n_3 = 25$.

at least in the range $0.3 < w/a < 0.9$ which is more useful.

D. Field beyond the optical waveguide

Propagation through free space beyond the waveguide of the emerged field considered in Figure 7 is shown in Figure 8.

We underline that beyond the waveguide, propagation through any optical path, composed by an apertured

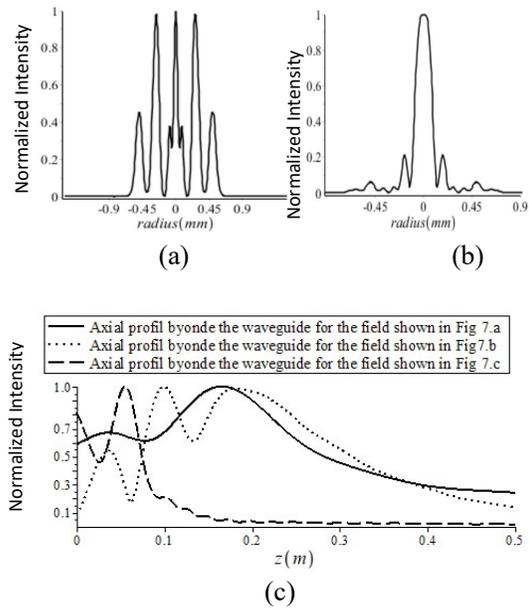


Fig. 8. Transverse distribution of the field shown in Figure 7 (b) at (a) $z = 15$ mm and (b) $z = 95$ mm beyond the optical waveguide, respectively, and (c) axial field distribution for the fields shown in Figure 7.

first-order optical system, can be achieved using the ABCD law and the generalized Gouy phase (GGP) shift [16–19].

III. CONCLUSION

In this paper, waveguide crossing has been solved by giving a simple rule for an efficient mode coupling at both ends of a cylindrical optical waveguide. Using the optimum number of modes given by the rule, crossing of the optical waveguide is achieved successfully regardless of the constraint of the ratio between the beam spot size and the waveguide radius.

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