

Numerical Analysis on the Equivalent Physical Temperature Estimation for Pyramidal Microwave Calibration Targets Based on Infrared Imagery

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Abstract – Pyramidal microwave calibration targets (MCT) are widely applied in on-orbit and pre-launch radiometric calibrations due to its compact size. However, it is well known that array-shaped MCT suffers from the temperature gradient at tips that leads to radiating brightness temperature (*BT*) bias. Therefore, it is vital to estimate the equivalent physical temperature of the MCT given the tip temperature gradient for the *BT* bias correction. In this work, the authors numerically investigate the T_{MCT} estimation based on possible temperature measurement techniques, specifically considering a combination configuration of platinum resistor temperature (PRT) detector at the metal base and the infrared camera detector for the tipbottom temperature difference. By considering the possible variation of coating material parameters and thermal measurement errors, it is possible to evaluate T_{MCT} estimation accuracy. Numerical results indicate that this temperature measurement configuration can lead to accurate T_{MCT} estimation at the level of 0.1 K ($1 - \sigma$). Factors that notably impact on the estimation accuracy are discussed. This investigation can be a direct reference for MCT *BT* correction applications in the pre-launch calibration process.

Index Terms – Brightness temperature estimation, microwave calibration target, pre-launch calibration.

I. INTRODUCTION

The microwave calibration target (MCT), specifically the array-shaped coated sharp pyramidal target, is widely applied in radiometric payloads [1–3]. The array-shaped MCT has the advantage of compactness in size and weight, which is beneficial for space-borne or air-borne applications [2–4], compared to other types

of MCT such as cavity MCT [5, 6]. However, the key problem of the array-shaped MCT in calibration application is that a temperature gradient exists at the tip region, and it leads to radiated brightness temperature (*BT*) bias.

As the function of the MCT is to provide an accurate referencing *BT*, efforts has been made on estimating and determining the radiation performance of the MCT. After a lot of investigations were performed upon the emissivity issue of the MCT, such as wide-band reflectivity spectrum properties [7], wideband high-emissivity design strategy [8, 9], and emissivity measurement methodology [10–14], research focus has been altered towards the overall *BT* performance considering temperature gradient at tips. In [15], Schröder et al. proposed a numerical framework to calculate the radiated *BT* from the arrayshaped MCT for space-borne application. In [16], Virone et al. calculated the radiated *BT* for a cold source in the vacuum-thermal calibration scenario. It is evidenced in numerous works that the temperature gradient becomes the primary source of *BT* bias in general cases over the residual reflectivity from the array-shaped MCT. In [17, 18], Jin et al. proposed a *BT*bias suppressed unit kernel design of array-shaped MCT based on the numerical framework in [15].

In [15–17], numerical frameworks are established for numerically estimating the *BT*-bias, which offers rich information and deep understanding of *BT*-bias properties in space-borne platform deployment and prelaunch thermal-vacuum (T/V) tests. However, the key concepts of those works are estimations based on numerical simulations as a reference for practical application. For practical pre-launch tests, it is desired that the *BT*-bias can be obtained based on measurable values other than fully based on simulations. In the earlier studies in which the temperature gradient phenomenon of the

array-shaped MCT was presented, infrared imagery was applied, showing the ability to capture the temperature difference between the MCT unit tips and unit valley bottoms [2, 19]. Based on the recent progress of BT modeling methodology, it is highly applicable to use infrared imagery to obtain temperature gradient information and to estimate BT -bias with numerically concluded electromagnetic information.

In this work, the authors investigate such a possibility. Specifically, the accuracy of the BT -bias estimation is investigated considering the uncertainty of MCT coating material properties, as well as the measurable temperature difference between tips and valley bottoms. The results can be a direct reference for BT -bias estimation in ground radiometer tests such as pre-launch T/V tests.

II. METHODOLOGY FOR T_{MCT} ESTIMATION

Due to the typical MCT structure of 2-D array of coated sharp pyramids, the dominating factor for BT bias is the temperature gradient in the coating layers in the height dimension (along the z -axis). Therefore, the simulation study for the T_{MCT} estimation, is focused on the temperature gradient in the height dimension.

A. Theoretical description for equivalent temperature estimation

In [15, 17], it is concluded that the directional radiated BT from the MCT can be calculated by a cross-integration of temperature and electromagnetic absorption distribution:

$$BT = (1 - r) \int_{CoatingLayer} T(\vec{r}) \tilde{A}(\vec{r}) dv + rBT_b, \quad (1)$$

where $T(\vec{r})$ is temperature distribution in the coating layer and $\tilde{A}(\vec{r})$ is normalized absorption rate distribution in the coating layer:

$$\int_{CoatingLayer} \tilde{A}(\vec{r}) dv = 1. \quad (2)$$

Here, r is total reflectivity of the MCT to normal incident electromagnetic waves and BT_b is the background brightness temperature and is considered as uniform in this investigation.

Based on the former studies, it is known that the most important temperature gradient phenomenon is the tip-bottom gradient (along z -direction), due to the difference of the bottom temperature T_{base} of MCT and the background BT_b . Equation (1) can be simplified as:

$$BT = (1 - r) \int_{CoatingLayer} T(z) \tilde{A}(z) dz + rBT_b. \quad (3)$$

Specifically, considering numerical investigations, $T(z)$ shall be calculated by thermal simulations, $\tilde{A}(z)$

and r shall be calculated by electromagnetic simulations. Equation (3) can be simplified as:

$$BT = eT_{MCT} + rBT_b, \quad (4)$$

where T_{MCT} can be defined as the equivalent physical temperature taking count of the temperature gradient in the coating layer:

$$T_{MCT} = \int_{CoatingLayer} T(z) \tilde{A}(z) dz. \quad (5)$$

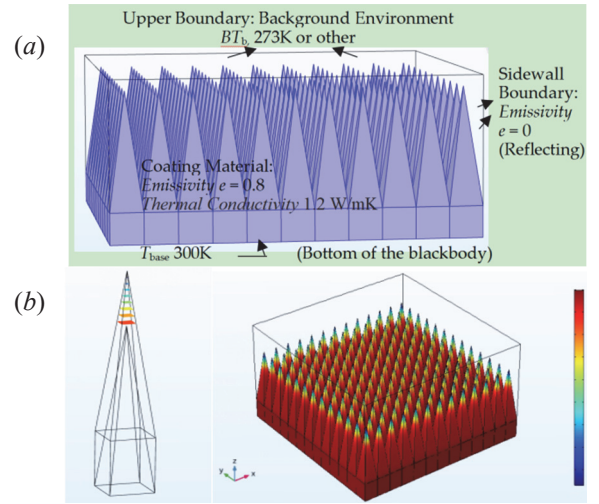


Fig. 1. Simulated temperature gradient of MCT array (in front of a colder background) showing notable temperature gradient at tips.

As the accuracy of radiated BT from the MCT is now a research focus for the microwave radiometer calibration, it was concluded that the difference between T_{MCT} and the ideal T_{base} is the key factor for the BT bias, as it generates the bias (ΔBT_1 in equation 6) notably larger than that by the reflectivity factor (ΔBT_2), especially at current stage that the emissivity of the MCT can be above 0.9995 even 0.9999 [4, 14]. Although T_{MCT} is a key factor for BT bias, unfortunately it cannot be directly measured based on current techniques. Therefore, it is important to develop estimation methodology according to the physical properties of impacting factors such as $T(z)$ and $\tilde{A}(z)$.

$$\Delta BT = e(T_{MCT} - T_{base}) + r(BT_b - T_{base}) = \Delta BT_1 + \Delta BT_2. \quad (6)$$

In former studies, it was found that T_{MCT} is proportional to the maximum temperature gradient at the tip ($\Delta T_{tip} = T_{tip} - T_{base}$) [15], as expressed in equation (7), and the corresponding proportionality factor $\alpha_{Gradient}$ can be numerically concluded. Also, it should be noted that, as the $\tilde{A}(z)$ is frequency-dependent, thus

so as the α_{Gradient} , then the T_{MCT} shall be frequency dependent. Further, based on the above acknowledgement, the authors view the temperature measurements realizable in cases of ground radiometer calibration, as shown in Fig. 2, can be a practical path to estimate the T_{MCT} with help of numerical concluded values. It is specifically suggested to measure the T_{base} using the platinum resistor temperature (PRT) detector (T_{PRT}), which is a widely-applied contact measurement technique with excellent precision. It is suggested to use an infrared imaging camera to capture the difference between the tip temperature and the bottom temperature (ΔT_{tip}) as the maximum temperature gradient. Based on this configuration, the aim of this study is to investigate and conclude the achievable accuracy of this potential methodology.

$$T_{\text{MCT}} = T_{\text{PRT}} + \alpha_{\text{Gradient}} \Delta T_{\text{tip}}. \quad (7)$$

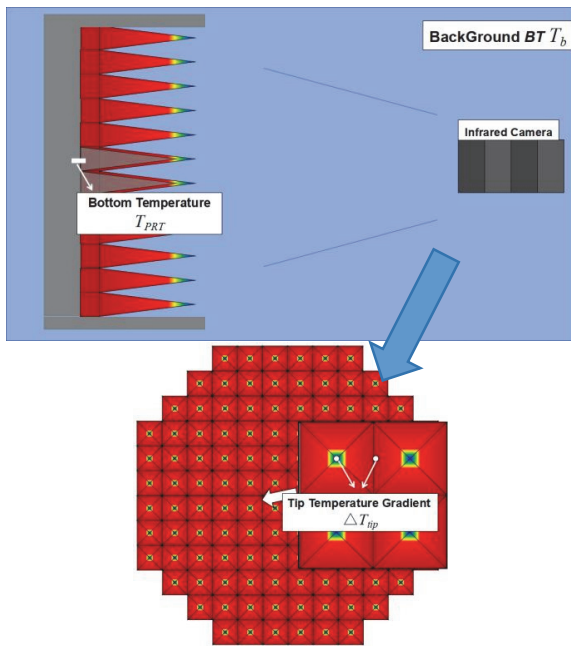


Fig. 2. Configuration for MCT temperature measurement.

It should be noted that estimating the T_{MCT} (equivalent physical temperature of the MCT) has been done in former studies based on comprehensive fullroutine numerical computations, resulting in the acknowledgement of important conclusions such as T_{MCT} is proportional (α_{Gradient}) to the maximum temperature gradient, and the proportional factor α_{Gradient} is frequency dependent. Based on that, the key aim of this work is to investigate the estimating accuracy for T_{MCT} based on measurable temperature values. This work is an update of the measurementnumerical routine from the stage of full-numerical estimation routine. The authors believe

such an update is an important reference for practical applications, showing how a good estimating accuracy can be achieved.

B. Discussion of thermal simulations

Before the T_{MCT} estimating methodology is investigated, it is important to review the thermal simulation for the temperature distribution, as shown in Fig. 1 (a). The temperature gradient in the coating layers is due to difference between the base temperature and background BT, and it can be simulated based on a finite element software package (such as COMSOL). It should be noted that clearly different input thermal parameters, such as the background BT or the emissivity of the coating layer, will lead to different temperature distribution results. However, it was found that, after normalization, the average temperature distribution results are nearly the same [20], as shown in Fig. 3. This may be because of the structure of the MCT unit, where the coating layer with low thermal conductivity and high emissivity is located above the highly thermal conducting metal kernel, and such a combination of parts with opposite thermal properties leads to a stable normalized temperature distribution in the coating layer. Based on this important phenomenon, it can be inferred that as one can obtain the maximum temperature difference between the tip and the bottom, one gets the overall temperature distribution and will be able to calculate T_{MCT} . This is the key reason that the infrared temperature camera is used for high-accuracy T_{MCT} estimations, as the most important information is the temperature difference between the tip and bottom, as a relative value, other than the absolute temperature value at each position. The accuracy of the absolute temperature value can be counted on the PRT measurement, which is found to be promising in practical experiences.

Before further investigations, it is also interesting to check the possible temperature gradient at the tip. As shown in Fig. 4, the tip-bottom temperature difference results in cases of different background BT plotted. Considering the applications in ground T/V tests, where 300 K hot MCT is in front of a colder background or an 80 K cold MCT is in front of a warmer background, the range of 30–160 K base-background temperature difference will leads to the range of 3–11 K base-tip temperature difference, showing a non-linear projection relationship. The considered upper-range of the base-tip temperature difference is 12 K, which is sufficient for the reference purpose.

C. Discussion of electromagnetic simulations

Electromagnetic simulations are conducted in the unit level to obtain the local absorption rate distribution

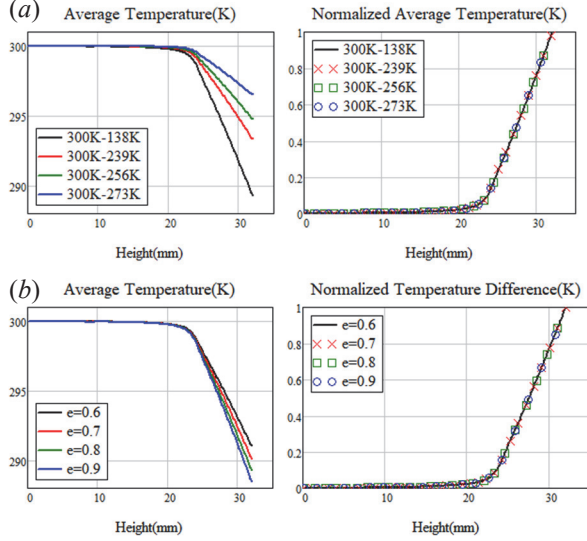


Fig. 3. Calculated $T(z)$ distributions considering different thermal parameters: (a) $e = 0.8$ and (b) temperature parameters of 300 K (MCT)-138K (Environment).

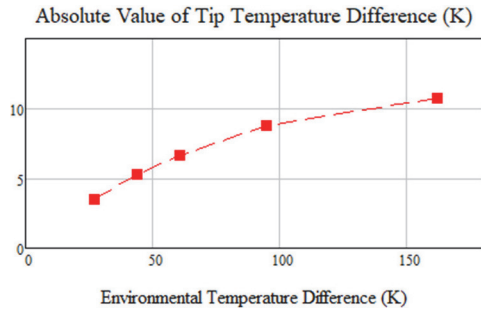


Fig. 4. Calculated tip-bottom temperature difference versus environmental temperature difference (K).

of $\tilde{A}(z)$. Specifically, based on the periodic boundary condition, one single unit is included in the computation domain of Finite Difference Time Domain (FDTD), where a plane-wave excitation is illuminating toward the unit along the z -direction and the total reflectivity r and local absorption rate $\tilde{A}(z)$ can be recorded [17, 18]. It should be noted that $\tilde{A}(z)$ is very hard to actually measure, and numerical calculation is necessary. The electromagnetic parameters of the coating layer material, such as the dielectric constant (ϵ), and the magnetic permeability (μ), will directly impact on the pattern of $\tilde{A}(z)$. In practical applications, the uncertainty of the (ϵ, μ) measurement, as well as the heterogeneity in the coating layer, will have negative impact on the accuracy of the calculated $\tilde{A}(z)$ in the process of T_{MCT} estimation. Therefore, before the overall investigation of T_{MCT} estimation accuracy can be performed, the uncertainty of $\tilde{A}(z)$ due to the uncertainty of (ϵ, μ) shall be investigated.

III. ACCURACY ASSESSMENT FOR T_{MCT} ESTIMATION

In this section, the numerical results are presented and discussed for the assessment of the T_{MCT} estimation accuracy.

A. Uncertainty of $\alpha_{Gradient}$ due to electromagnetic parameter variation

Uncertainty of the (ϵ, μ) of the coating layer material will lead to variation of the $\tilde{A}(z)$ then the counted $\alpha_{Gradient}$. For this issue, the authors set up the randomly distributed (ϵ, μ) at different levels of standard derivation error, and performed 10 rounds of simulations to count the resulted $\alpha_{Gradient}$. For each random sample, $\tilde{A}(z)$ is calculated and integrated with $T(z)$ for T_{MCT} , then $c\alpha_{Gradient}$ as in equation (8). The authors considered other factors such as different frequency and different unit size, as in Tables 1 and 2. The considered frequencies, 10.65 GHz, 18.7 GHz, 36.5 GHz, 89 GHz, 183.31 GHz, are the most popular frequency channels applied in the radiometer payloads, and the performance of the MCT at those frequency points are favored by radiometer developers:

$$\alpha_{Gradient} = \frac{T_{MCT} - T_{base}}{\Delta T_{tip}}, \quad (8)$$

Table 1: Geometry parameters for MCT

Case	Period	Height to Period	Coating Thickness
Case 1	6.0 mm	□4 : 1	1.0 mm
Case 2	7.5 mm	□4 : 1	1.0 mm

Table 2: Considered coating material electromagnetic parameters at different frequencies

Frequencies	Dielectric Constant	Magnetic Permeability
10.65 GHz	10.284 – 0.185j	1.511 – 0.872j
18.7 GHz	10.253 – 0.32j	1.203 – 0.594j
36.5 GHz	10.133 – 0.588j	0.736 – 0.205j
89 GHz	9.604 – 1.06j	0.956 – 8.165e – 3j
183.31 GHz	8.752 – 1.161j	0.99 – 8.563e – 4j

Note: The parameters are from the MF114 description based on the measured data in [20].

Based on the initial parameters, the calculated results are presented in Figs. 5–7. In Figs. 5 and 6, the calculated local absorption rates $\tilde{A}(z)$ are presented in cases of different unit size and at different frequencies. It is clear that $\tilde{A}(z)$ is highly frequency dependent. A higher absorption rate in the tip region will lead to a higher $\alpha_{Gradient}$. At 36.5 GHz and 89 GHz the counted

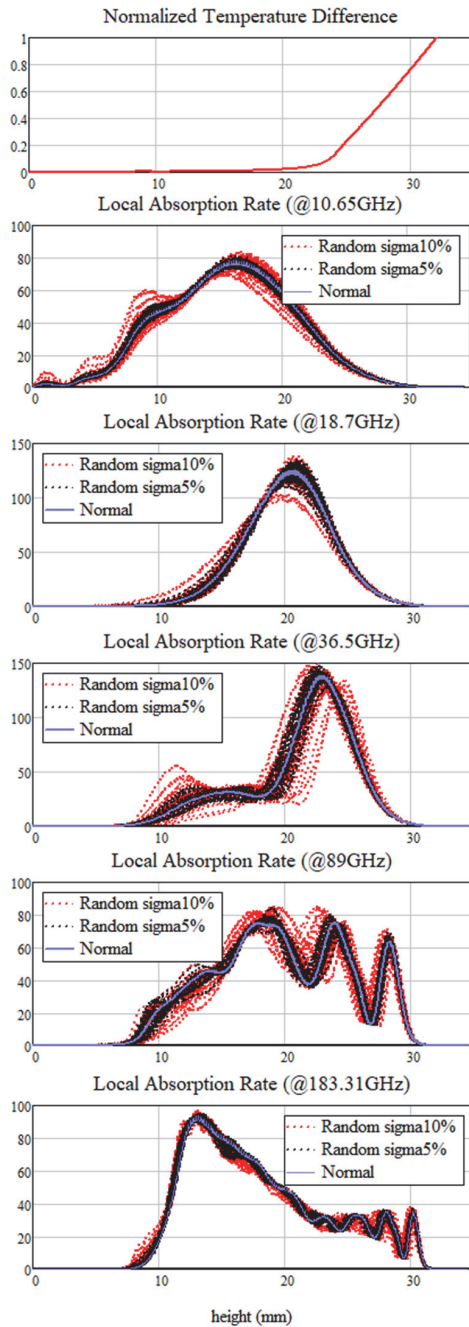


Fig. 5. Calculated local absorption distribution results in cases of different frequency, considering random samples, and referenced to the average temperature distribution in the coating layer. Unit period size $p = 6.0$ mm, considering random samples with 5% std in (ϵ, μ) and with 10% std in (ϵ, μ) .

α_{Gradient} will be relatively higher. In Figs. 5 and 6, the $\tilde{A}(z)$ curves for random samples are also plotted. One can intuitively observe the influence of the uncertainty of the material (ϵ, μ) , upon the $\tilde{A}(z)$ which directly affects the α_{Gradient} .

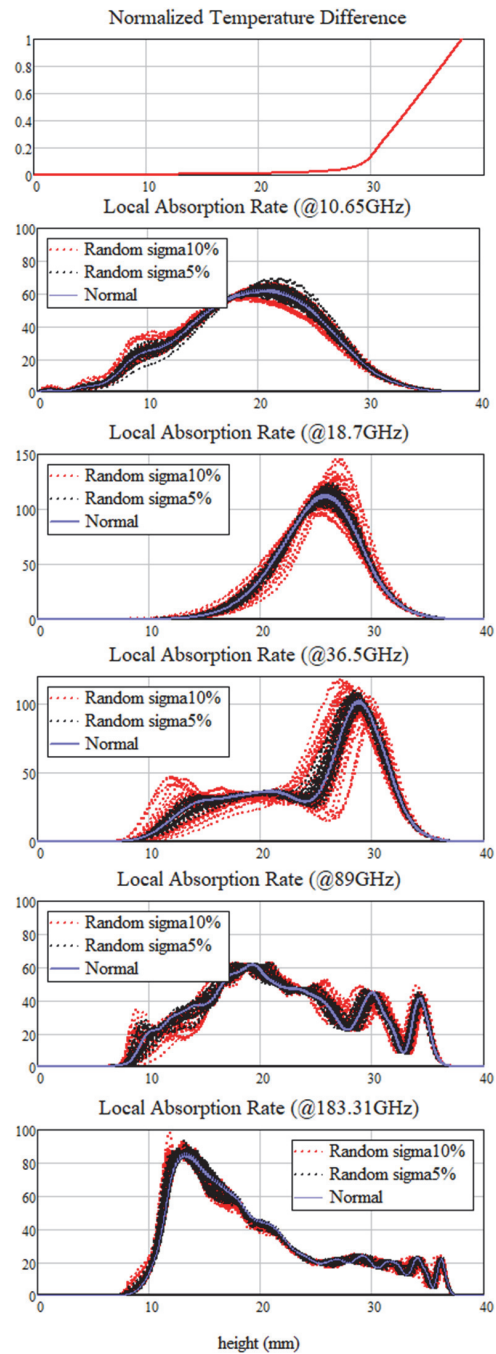


Fig. 6. Calculated local absorption distribution results in cases of different frequency, considering random samples, and referenced to the average temperature distribution in the coating layer. unit period size $p = 7.5$ mm, considering random samples with 5% std in (ϵ, μ) and with 10% std in (ϵ, μ) .

In Fig. 7, the calculated α_{Gradient} and counted standard deviation σ (std for short) are plotted and compared. The data set contains the comparison between the two cases with different unit size, $p = 6.0$ mm and

$p = 7.5$ mm. In Figs. 7 (a) and (b), one can see that the calculated α_{Gradient} in case of $p = 6.0$ mm is clearly larger than that of $p = 7.5$ mm at every frequency. When $p = 6.0$ mm, the highest α_{Gradient} is found at 89 GHz. When $p = 7.5$ mm, the highest α_{Gradient} is found at 36.5 GHz. Thus, the α_{Gradient} -frequency relationship is variable versus the geometry size. For std, the counted value for α_{Gradient} when $p = 6.0$ mm is generally larger than when $p = 7.5$ mm, but not at every frequency point. When $p = 6.0$ mm and $p = 7.5$ mm, the largest counted std is found at 36.5 GHz.

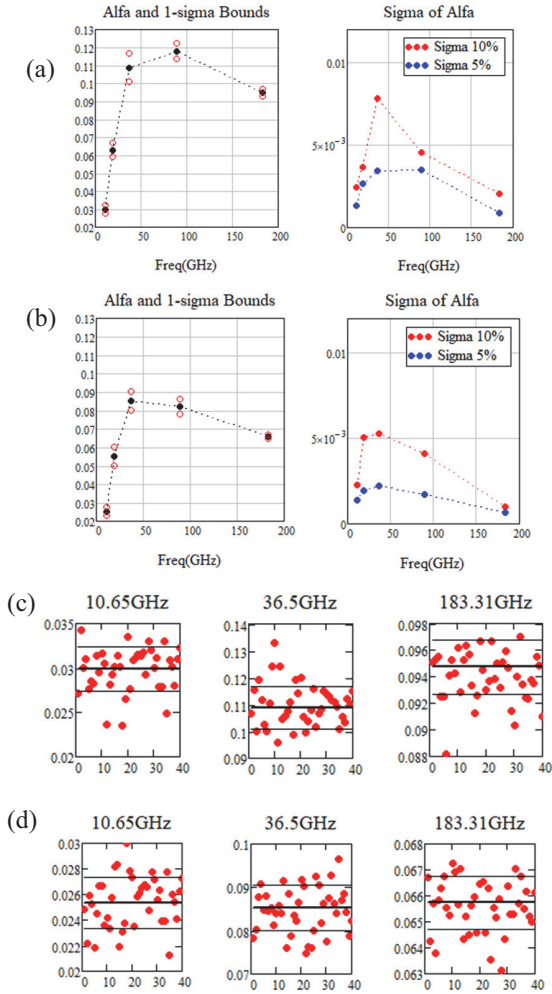


Fig. 7. Counted α_{Gradient} values based on 40 random samples: (a) $p = 6.0$ mm, counted α_{Gradient} at different frequencies, and counted std values in cases of (ϵ, μ) std at 5% and 10%; (b) $p = 7.5$ mm, counted α_{Gradient} at different frequencies, and counted std values in cases of (ϵ, μ) std at 5% and 10%; (c) $p = 6.0$ mm, counted α_{Gradient} at each random sample at 10.65 GHz, 36.5 GHz and 183.31 GHz ; (d) $p = 7.5$ mm, counted α_{Gradient} at each random sample at 10.65 GHz, 36.5 GHz and 183.31 GHz.

In these set of results, another important fact is that a smaller (ϵ, μ) std leads notably smaller α_{Gradient} std, which can be anticipated. The comparison between different (ϵ, μ) std level is for further investigation upon the T_{MCT} , showing the importance of the (ϵ, μ) error control. For obtaining the statistical std information, 40 random samples are calculated for each frequency and each (ϵ, μ) sigma configuration. The calculations were implemented on one workstation with EPYC 32-core CPU and took more than 10 minutes for each sample. In running each sample, most of the computation time was consumed on performing Fourier transform at each cell for storing local fields at the specific frequency, which were then used for obtaining the local absorption distribution.

B. Uncertainty evaluation for T_{MCT} estimation

Based on the counted α_{Gradient} and its std results, it is further possible to evaluate T_{MCT} estimation accuracy. As discussed in the previous section, there are several obstructions for accurate thermal simulation. In this case, one may not expect an accurate BT bias estimation based on simulated results. However, it is important that the relative temperature distribution across the pyramidal unit is fixed, and it can be inferred that the key issue of temperature measurement is to obtain the maximum value of thermal gradient, which is the bottom-tip difference ΔT_{tip} .

Considering the bias of the T_{MCT} , as ΔT_{MCT} described in (9), then it is possible to estimate the overall σ (std) based on Monte Carlo evaluation:

$$\Delta T_{\text{MCT}} = T_{\text{PRT}} + \alpha_{\text{Gradient}} \Delta T_{\text{tip}} - T_{\text{base}}. \quad (9)$$

Specifically, the T_{base} is considered as the absolute value, and:

- (1) T_{PRT} is the measurable base temperature by the PRT detector, a σ_{TPRT} at the level of 0.03 K is considered. Actually, the PRT temperature measurement has been proven to be accurate and reliable in practical MCT application.
- (2) ΔT_{tip} is the temperature difference between the unit tip and the bottom, measurable by the infrared camera detection. Actually, the sensitivity for temperature magnitude measurable by the infrared camera can be at the level of 0.1 K. Also, accurate infrared camera detection will rely on accurate estimation of surface emissivity, therefore, the $\sigma_{\Delta T_{\text{tip}}}$ shall be notably larger than 0.1 K. In this work, $\sigma_{\Delta T_{\text{tip}}}$ is set to be 0.5 K.
- (3) $\sigma_{\alpha_{\text{Gradient}}}$ is the std for the calculated α_{Gradient} because of the uncertainty of the coating material electromagnetic parameters, which is frequency dependent and calculated in the former section.

With the above factors considered, the counted T_{MCT} estimation accuracy results are presented and discussed based on Monte Carlo analysis.

In Fig. 8, a set of estimated T_{MCT} bias results are plotted at specific parameters considering $\Delta T_{tip} = 10$ K. As can be observed, T_{MCT} bias varies versus frequency with a considerable large range. The T_{MCT} bias is generally larger in magnitude around the frequency $f_0 = c_0/p$, than in high frequency range (f is notably larger than f_0). T_{MCT} bias gets smaller as frequency f gets higher. Considering the comparison of unit period, a larger unit with the same coating parameter and height to period ratio generally leads to smaller T_{MCT} bias.

Considering the above uncertainty factors, the std σ results for the estimated T_{MCT} bias are also plotted. Several sets of comparisons are plotted in Figs. 9–11.

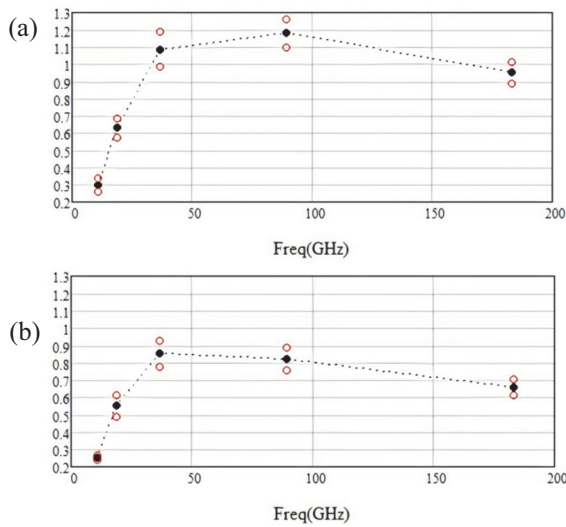


Fig. 8. Estimated T_{MCT} bias (absolute value of $T_{MCT} - T_{base}$) and $1 - \sigma$ std at different frequencies: (a) $p = 6.0$ mm and (b) $p = 7.5$ mm. $\sigma_\epsilon = 0.1\epsilon$, $\sigma_\mu = 0.1\mu$, $\Delta T_{tip} = 10$ K, $\sigma\Delta T_{tip} = 0.5$ K.

In Fig. 9, counted T_{MCT} bias std at different frequencies are compared. One can see the accuracy improvement contributed by achieving higher (ϵ, μ) determination accuracy. An interesting and important phenomenon can be observed in Fig. 9 (a). As the maximum tip temperature gradient is rising, the std of T_{MCT} estimation also gets larger, especially at center frequencies (36.5 GHz). This trend is less notable in Fig. 8 (a). This shows the importance of the coating material parameter determination accuracy, especially at central frequencies ($f_0 = c_0/p$). A sufficiently low uncertainty in the coating layer material parameter will contribute to a stably low std in the T_{MCT} estimation, even when the environmental temperature difference is quite large ($T_{base} - BT_b$).

In Fig. 10, similar std results of the T_{MCT} bias estimation are compared but considering a larger unit size $p = 7.5$ mm. Overall data trends are similar to those in Fig. 9. One can see that the std for the T_{MCT} estimation are notably smaller than that in Fig. 8. The key reason for this fact is that the $\alpha_{Gradient}$ in case of a larger unit is smaller than that in case of a smaller unit, given the same coating thickness and height to period ratio. This set of comparison shows that a larger unitsize MCT design generally leads better T_{MCT} estimation accuracy.

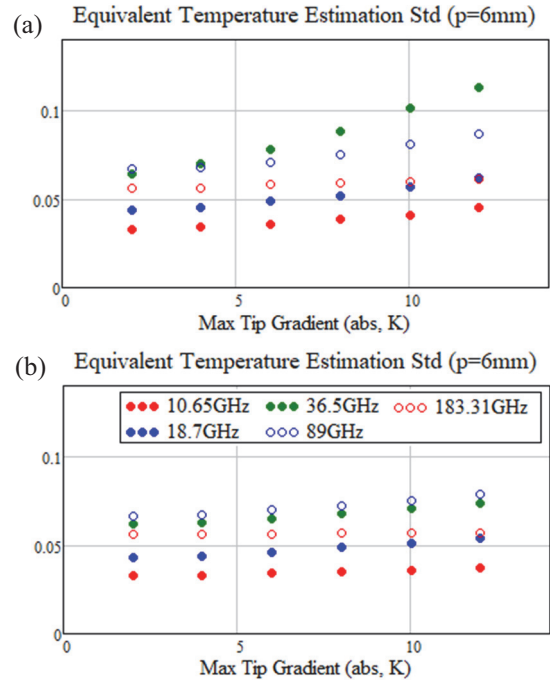


Fig. 9. Counted T_{MCT} bias std at different frequencies: (a) $\sigma_\epsilon = 0.1\epsilon$, $\sigma_\mu = 0.1\mu$ and (b) $\sigma_\epsilon = 0.05\epsilon$, $\sigma_\mu = 0.05\mu$. $p = 6.0$ mm, $\sigma\Delta T_{tip} = 0.5$ K.

In Figs. 9 and 10, most of the std results are below 0.1 K, which is a considerably small value sufficient for practical BT correction. It is also important to recall the numerical analysis for the BT bias estimation in V/T test reported in [16], which gives many important insights into the BT radiation process. However, based on a full-numerical estimation approach, the counted BT bias estimation std is at the level of 0.22 K (maximum) for the MCT $p = 8.6$ mm. Considering those results, the key motivation of this work shall be concluded as: shorten the error analysis chain by considering measurable temperature values over simulated results will lead to the BT bias estimation more practical and more reliable.

Consider a higher rise in the tip-bottom temperature difference measurement by infrared camera detection, which is highly possible in practical applications, specifically $\sigma\Delta T_{tip} = 1$ K. The corresponding results are given

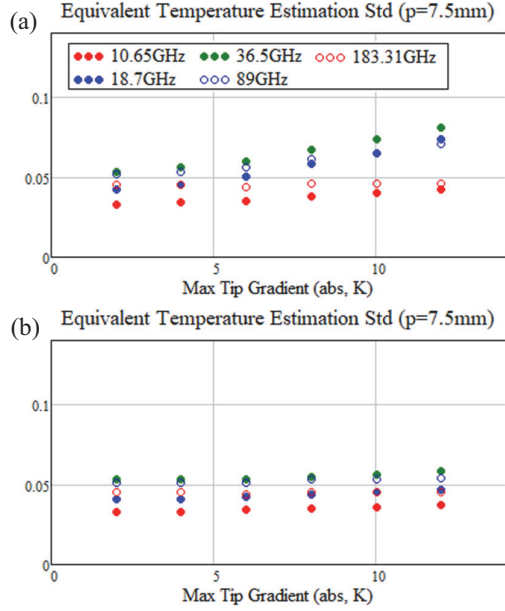


Fig. 10. Counted T_{MCT} bias std at different frequencies: (a) $\sigma_\epsilon = 0.1\epsilon, \sigma_\mu = 0.1\mu$ and (b) $\sigma_\epsilon = 0.05\epsilon, \sigma_\mu = 0.05\mu, p = 7.5 \text{ mm}, \sigma\Delta T_{tip} = 0.5 \text{ K}$.

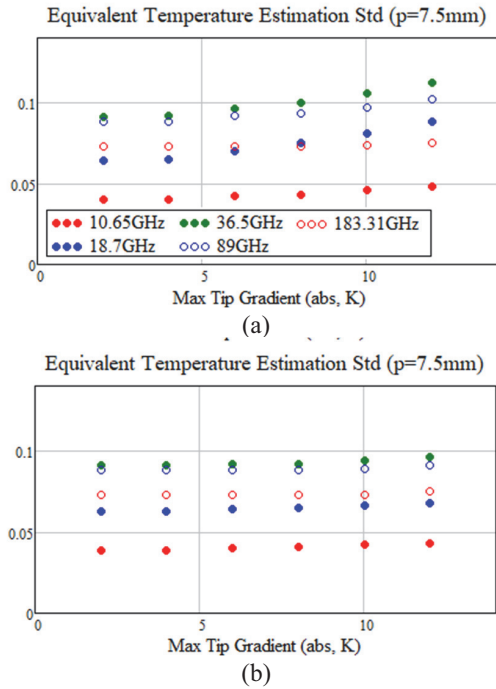


Fig. 11. Counted T_{MCT} bias std at different frequencies: (a) $\sigma_\epsilon = 0.1\epsilon, \sigma_\mu = 0.1\mu$ and (b) $\sigma_\epsilon = 0.05\epsilon, \sigma_\mu = 0.05\mu, p = 7.5 \text{ mm}, \sigma\Delta T_{tip} = 1 \text{ K}$.

in Fig. 11. By achieving a sufficient (ϵ, μ) accuracy, $\sigma_\epsilon = 0.05\epsilon, \sigma_\mu = 0.05\mu$, the overall std of the T_{MCT} estimation can be below 0.1 K. This set of results indicates that T_{MCT} estimation based on PRT measurement

at bottom and infrared camera detection for tip-bottom temperature difference is generally robust and can afford the relatively low accuracy for the ΔT_{tip} measurement.

IV. CONCLUSIONS

In this work, the authors investigate the equivalent physical temperature (T_{MCT}) estimation methodology for the MCT in ground-based radiometer calibration applications. It is well known that the difference between the actually T_{MCT} and the nominal temperature is the key factor for the BT bias. Accurate estimation of T_{MCT} is significant for practical BT bias correction. This work performs error analysis for T_{MCT} estimation, for a highly possible temperature measurement strategy combing PRT and infrared camera detection. It is shown in the numerical results that such a routine is possible to determine T_{MCT} at an accuracy level of $1 - \sigma$ std 0.1 K, which is important for practical calibration application. This routine performs better in accuracy on the MCT with a larger unit size over that with a smaller unit size. Ensuring sufficient low uncertainty in the electromagnetic parameters of the coating layer is important to retain the overall estimation accuracy and make such a routine more tolerable to the infrared detection error. Numerical results show the possibility of the temperature measurement supported BT bias estimation and provide a corresponding reference for error analysis.

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